

Chapter III

Light-matter interaction processes

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Outline

- 1) Light-matter interaction Hamiltonian
- 2) Emission and absorption
- 3) Light scattering

1) Light-matter interaction Hamiltonian

For one particle:
$$\hat{H} = \frac{1}{2m} \left[\hat{\vec{P}} - q\vec{A}(\hat{\vec{R}}) \right]^2 + qU(\hat{\vec{R}})$$

$$\hat{H} = \frac{\hat{\vec{P}}^2}{2m} - \frac{q}{2m} \underbrace{\left[\hat{\vec{P}} \cdot \vec{A}(\hat{\vec{R}}) + \vec{A}(\hat{\vec{R}}) \cdot \hat{\vec{P}} \right]}_{[\hat{\vec{R}}, \hat{\vec{P}}] \neq 0} + qU(\hat{\vec{R}}) + \frac{q^2 \vec{A}^2(\hat{\vec{R}})}{2m}$$

Many-particle:

$$\begin{aligned} \hat{H} &= \sum_i \left\{ \frac{\hat{\vec{P}}_i^2}{2m} - \frac{q}{2m} \left[\hat{\vec{P}}_i \cdot \vec{A}(\hat{\vec{R}}_i) + \vec{A}(\hat{\vec{R}}_i) \cdot \hat{\vec{P}}_i \right] + qU(\hat{\vec{R}}_i) + \frac{q^2 \vec{A}^2(\hat{\vec{R}}_i)}{2m} \right\} \\ &= \hat{H}_0 + \hat{H}_{int} . \end{aligned}$$

(same masse m and same charge q)

1) Light-matter interaction Hamiltonian

Paramagnetic current operator:

$$\hat{\vec{J}}(\vec{r}) \equiv \frac{1}{2m} \sum_i \left\{ \hat{\vec{P}}_i \delta(\vec{r} - \hat{\vec{R}}_i) + \delta(\vec{r} - \hat{\vec{R}}_i) \hat{\vec{P}}_i \right\}$$

Density operator: $\hat{n}(\vec{r}) \equiv \sum_i \delta(\vec{r} - \hat{\vec{R}}_i)$

$$\hat{H}_{int} = \int d\vec{r} \left\{ -q \hat{\vec{J}}(\vec{r}) \cdot \vec{A}(\vec{r}) + \frac{q^2}{2m} \hat{n}(\vec{r}) \vec{A}^2(\vec{r}) + q \hat{n}(\vec{r}) U(\vec{r}) \right\}$$

$$\hat{H}_{int} = \hat{H}_1 + \hat{H}_2 + \hat{H}_3$$

two choices

$\vec{A}(\vec{r})$: Classical description of the field

$\hat{\vec{A}}(\vec{r})$: Quantum description of the field

$$\hat{H}_1 = -q \int d\vec{r} \hat{\vec{J}}(\vec{r}) \cdot \vec{A}(\vec{r}) ,$$

$$\hat{H}_2 = \frac{q^2}{2m} \int d\vec{r} \hat{n}(\vec{r}) \vec{A}^2(\vec{r}) ,$$

$$\hat{H}_3 = q \int d\vec{r} \hat{n}(\vec{r}) U(\vec{r}) .$$

1) Light-matter interaction Hamiltonian

Using $\hat{\vec{A}}(\vec{r}) = \frac{1}{\sqrt{V}} \sum_{i, \vec{k}} \sqrt{\frac{\hbar}{2\omega_k \epsilon_0}} \left\{ \hat{a}_{i, \vec{k}} \hat{\epsilon}_{i, \vec{k}} e^{i\vec{k} \cdot \vec{r}} + \hat{a}_{i, \vec{k}}^+ \hat{\epsilon}_{i, \vec{k}}^* e^{-i\vec{k} \cdot \vec{r}} \right\}$

$$\Rightarrow \hat{H}_1 = \frac{e}{\sqrt{V}} \sum_{i, \vec{k}} \sqrt{\frac{\hbar}{2\omega_k \epsilon_0}} \left\{ \hat{a}_{i, \vec{k}} \hat{\epsilon}_{i, \vec{k}} \cdot \hat{\vec{J}}_{-\vec{k}} + \hat{a}_{i, \vec{k}}^+ \hat{\epsilon}_{i, \vec{k}}^* \cdot \hat{\vec{J}}_{\vec{k}} \right\}$$

with $\hat{\vec{J}}_{\vec{k}} = \int d\vec{r} \hat{\vec{J}}(\vec{r}) e^{-i\vec{k} \cdot \vec{r}}$

2) Emission and absorption

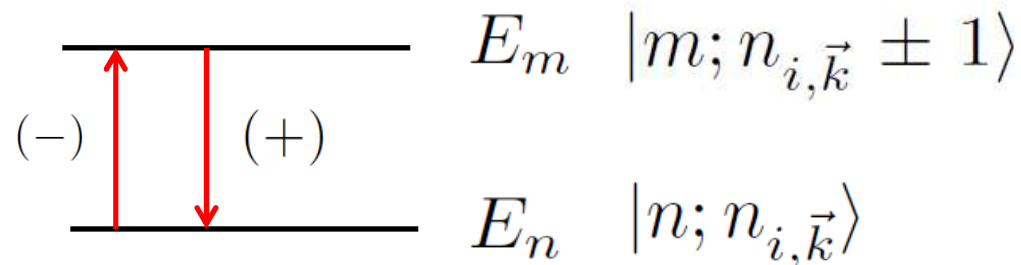
2.1) Emission and absorption rates

**Fermi
Golden
Rule**

$$\Gamma = \frac{2\pi}{\hbar} \left| \langle m; n_{i,\vec{k}} \pm 1 | \hat{H}_1 | n; n_{i,\vec{k}} \rangle \right|^2 \rho_f^\pm$$

(FGR)

First-order
perturbation theory



E_n Energies of the atomic, molecular, band...states

ρ_f Density of final states $\omega = c|\vec{k}| \equiv \omega_k$

(-) absorption $\rho_f^- = \delta(E_m - E_n - \hbar\omega)$

(+) emission $\rho_f^+ = \int \frac{d\vec{r}d\vec{p}}{(2\pi\hbar)^3} \delta(E_m - E_n + \hbar\omega)$

2) Emission and absorption

Note that the states of radiation can also contain **any number of photons** in other modes, but here we focus on a process in which **a single photon** is emitted or absorbed, so that these other photons are **passive spectators** who have no influence on the amplitude of the process.

The hamiltonian \hat{H}_I is a sum over all the *em* modes but only one term of this sum contributes to the amplitude: the one associated to the mode $n_{i,\vec{k}}$

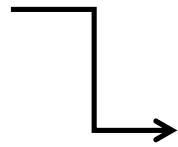
\hat{H}_I is a product of an operator acting on the Hilbert space of the field and an operator acting on the Hilbert space of the material system

$$\hat{H}_1 = \frac{e}{\sqrt{V}} \sum_{i,\vec{k}} \sqrt{\frac{\hbar}{2\omega_k \epsilon_0}} \left\{ \hat{a}_{i,\vec{k}} \hat{\epsilon}_{i,\vec{k}} \cdot \hat{\vec{J}}_{-\vec{k}} + \hat{a}_{i,\vec{k}}^+ \hat{\epsilon}_{i,\vec{k}}^* \cdot \hat{\vec{J}}_{\vec{k}} \right\}$$

↑ ↑
field matter

2) Emission and absorption

$$\Gamma = \frac{2\pi}{\hbar} \left| \langle m; n_{i,\vec{k}} \pm 1 | \hat{H}_1 | n; n_{i,\vec{k}} \rangle \right|^2 \rho_f^\pm$$

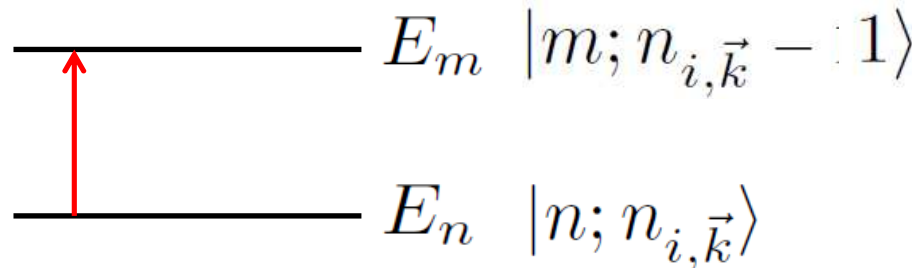


$$\begin{aligned} \langle n_{i,\vec{k}} + 1 | \hat{a}_{i,\vec{k}}^+ | n_{i,\vec{k}} \rangle &= \sqrt{n_{i,\vec{k}} + 1} \\ \langle n_{i,\vec{k}} - 1 | \hat{a}_{i,\vec{k}} | n_{i,\vec{k}} \rangle &= \sqrt{n_{i,\vec{k}}} \end{aligned}$$

electromagnetic part

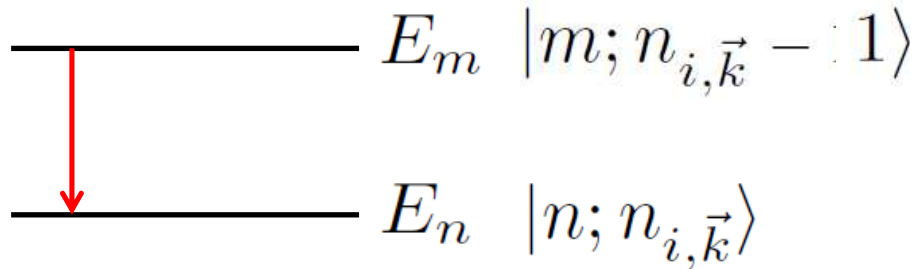
absorption

$$\Gamma_{abs.} = \frac{\pi e^2}{V \omega \epsilon_0} n_{i,\vec{k}} \left| \langle m | \hat{\vec{J}}_{-\vec{k}} \cdot \hat{\vec{\epsilon}}_{i,\vec{k}} | n \rangle \right|^2 \rho_f^-$$



2) Emission and absorption

emission



$$\Gamma_{em.} = \frac{\pi e^2}{V \omega \epsilon_0} (n_{i,\vec{k}} + 1) \left| \langle m | \hat{J}_{\vec{k}} \cdot \hat{\epsilon}_{i,\vec{k}}^* | n \rangle \right|^2 \rho_f^+$$

Most notable in these formulas is the presence of factors n and $n+1$ derived from what the photon is a boson and are responsible for the **emission** and **absorption** of radiation **induced**.

The more photons in the initial state, the greater the process of emission or absorption is likely.

This is the principle of the **laser** effect.

2) Emission and absorption

$$\Gamma_{em.} = \frac{\pi e^2}{V\omega\epsilon_0} (n_{i,\vec{k}} + 1) \left| \langle m | \hat{\vec{J}}_{\vec{k}} \cdot \hat{\vec{\epsilon}}_{i,\vec{k}}^* | n \rangle \right|^2 \rho_f^+$$

It is also remarkable that the emission of a photon can take place even in the absence of initial photon ($n = 0$), it means in the apparent absence of external stimulation.

Spontaneous emission

$$\Gamma_{em.} = \frac{\pi e^2}{V\omega\epsilon_0} \left| \langle m | \hat{\vec{J}}_{\vec{k}} \cdot \hat{\vec{\epsilon}}_{i,\vec{k}}^* | n \rangle \right|^2 \rho_f^+$$

It corresponds to the classical radiation of an accelerated charge.

